

# 高阶非线性海浪波面斜率的联合 概率统计分布\*

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**摘要** 基于非线性随机海浪模型,对各向异性波面斜率统计分布进行了理论研究。结果表明:在高阶近似下,波面斜率统计分布为截断的 Gram-Charlier级数,如果忽略波-波相互作用的非线性影响,则 Cram-Charlier级数分布蜕化为正态分布。

**关键词** 非线性统计分布 波面斜率 耦合矩

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海浪波面斜率非线性统计分布是深入研究海面和近海面小尺度海-气相互作用过程的重要基础,特别在海面遥感测量技术和海洋工程中得到广泛应用。迄今为止,对这一领域的研究已经取得了一系列重要成果(Jackson, 1979; Plant, 1986; Weissman, 1990; Resio, *et al*, 1991)。张书文等(1999)基于波面各相同性的假设,利用直接求耦合矩的方法,导出了具有对称性的波面斜率统计分布。本文将此项工作推广至各向异性的波面情形。

## 1 耦合矩的计算

依据 Longuet-Higgins 的非线性随机海浪模型,海浪波面高度  $\zeta$  表示为:

$$\zeta = \zeta_1 + \zeta_2 + \zeta_3 + \zeta_4 + \dots \quad (1)$$

由(1)式对位置坐标  $x, y$  的求偏导数,可得波面斜率  $\zeta_x, \zeta_y$  的表达式为:

$$\zeta_x = \zeta_{1x} + \zeta_{2x} + \zeta_{3x} + \zeta_{4x} + \dots \quad (2)$$

$$\zeta_y = \zeta_{1y} + \zeta_{2y} + \zeta_{3y} + \zeta_{4y} + \dots \quad (3)$$

现设  $\zeta_{1x}, \zeta_{1y}$  为一量级为  $\delta$  的小量,则  $\zeta_{nx}, \zeta_{ny}$  的量级应为  $\delta^n$ ,对于  $k$  近似下的波面斜率,当计算其  $n$  次方幂时,至少应精确至  $\delta^{n+k+1}$  的量级,在四阶近似下则有:

$$\zeta_x^n \zeta_y^m = \zeta_{1x}^n \zeta_{1y}^m + n \zeta_{1x}^{n-1} \zeta_{2x} \zeta_{1y}^m + m \zeta_{1x}^n \zeta_{1y}^{m-1} \zeta_{2y} + n \zeta_{1x}^{n-1} \zeta_{3x} \zeta_{1y}^m + m \zeta_{1x}^n \zeta_{1y}^{m-1} \zeta_{3y}$$

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$$\begin{aligned}
& + nm\zeta_{1x}^{n-1}\zeta_{2x}\zeta_{1y}^{m-1}\zeta_{2y} + \frac{n(n-1)}{2}\zeta_{1x}^{n-2}\zeta_{2x}^2\zeta_{1y}^m + \frac{m(m-1)}{2}\zeta_{1x}^n\zeta_{1y}^{m-2}\zeta_{2y}^2 + \\
& n\zeta_{1x}^{n-1}\zeta_{4x}\zeta_{1y}^m + m\zeta_{1x}^n\zeta_{1y}^{m-1}\zeta_{4y} + nm\zeta_{1x}^{n-1}\zeta_{2x}\zeta_{1y}^{m-1}\zeta_{3y} + nm\zeta_{1x}^{n-1}\zeta_{3x}\zeta_{1y}^{m-1}\zeta_{2y} + \\
& n(n-1)\zeta_{1x}^{n-2}\zeta_{2x}\zeta_{3x}\zeta_{1y}^m + m(m-1)\zeta_{1x}^n\zeta_{1y}^{m-2}\zeta_{2y}\zeta_{3y} + \frac{nm(n-1)}{2}\zeta_{1x}^{n-2}\zeta_{2x}^2\zeta_{1y}^{m-1}\zeta_{2y} \\
& + \frac{m(m-1)(m-2)}{6}\zeta_{1x}^n\zeta_{1y}^{m-3}\zeta_{2y}^3 + \frac{nm(m-1)}{2}\zeta_{1x}^{n-1}\zeta_{2x}\zeta_{1y}^{m-2}\zeta_{2y}^2 + \frac{n(n-1)(n-2)}{6}\zeta_{1x}^{n-3}\zeta_{2x}^3\zeta_{1y}^m
\end{aligned} \tag{4}$$

式中  $m, n$  为正整数。选取海面风的作用方向沿  $x$  方向,应用直接求矩的方法(孙孚等, 1994),可以导出

$$\begin{aligned}
\langle \zeta_x^{2N}\zeta_y^{2M} \rangle &= [(2N)!(2M)! / N!M!]r^N S^M [1 + a_{20}N/r + a_{02}M/S + \\
& a_{22}NM/rS + a_{40}N(N-1)/r^2 + a_{04}M(M-1)/S^2]
\end{aligned} \tag{5}$$

$$\langle \zeta_x^{2N}\zeta_y^{2M+1} \rangle = [(2N)!(2M+1)! / N!M!]r^N S^M [a_{21}N/r + a_{03}M/S] \tag{6}$$

$$\langle \zeta_x^{2N+1}\zeta_y^{2M} \rangle = [(2N+1)!(2M)! / N!M!]r^N S^M [a_{12}M/S + a_{30}N/r] \tag{7}$$

$$\langle \zeta_x^{2N+1}\zeta_y^{2M+1} \rangle = [(2N+1)!(2M+1)! / N!M!]r^N S^M [a_{11} + a_{31}N/r + a_{13}M/S] \tag{8}$$

式中  $N, M$  为正整数,  $r = \frac{1}{2} \int u^2 s(\vec{k}) d\vec{k}$ ,  $s = \frac{1}{2} \int v^2 s(\vec{k}) d\vec{k}$ ,  $s(\vec{k})$  为线性意义下的海浪波数谱, 相关系数  $a_{11}, a_{20}, a_{02}, a_{22}, a_{31}, a_{13}, a_{40}, a_{04}, a_{30}, a_{03}, a_{21}, a_{12}$  可由(9)式给出:

$$\begin{aligned}
a_{11} &= \langle B_1 C_1 \rangle + \langle B_2 C_2 \rangle + \langle B_3 C_3 \rangle + \langle B_1 C_3 \rangle; a_{20} = \langle B_1 B_3 \rangle + \langle B_2^2 \rangle / 2; a_{02} = \langle C_1 C_3 \rangle + \langle C_2^2 \rangle / 2; \\
a_{22} &= 4\langle B_1 B_2 C_1 C_2 \rangle + \langle B_2^2 C_1^2 \rangle + \langle B_1^2 C_2^2 \rangle + 2\langle B_1^2 C_2 C_3 \rangle + 2\langle B_1 B_3 C_1^2 \rangle; \\
a_{31} &= \langle B_1^3 C_3 \rangle + \langle B_1^2 B_2 C_2 \rangle / 2 + 3\langle B_1 B_2^2 C_1 \rangle; a_{13} = \langle B_3 C_1^3 \rangle + \langle B_2 C_1^2 C_2 \rangle / 2 + 3\langle B_1 C_1 C_2^2 \rangle \\
a_{40} &= 6\langle B_1^2 B_2^2 \rangle + 4\langle B_1^3 B_3 \rangle; a_{04} = 6\langle C_1^2 C_2^2 \rangle + 4\langle C_1^3 C_3 \rangle; a_{30} = 6\langle B_1 B_2 B_3 \rangle + 3\langle B_1^2 B_4 \rangle; \tag{9} \\
a_{03} &= 6\langle C_1 C_2 C_3 \rangle + 3\langle C_1^2 C_4 \rangle; a_{21} = \langle B_1^2 C_4 \rangle + 2\langle B_1 B_2 C_3 \rangle + 2\langle B_1 B_3 C_2 \rangle + 2\langle B_1 B_4 C_1 \rangle + 2\langle B_2 B_3 C_1 \rangle \\
a_{12} &= \langle B_4 C_1^2 \rangle + 2\langle B_3 C_1 C_2 \rangle + 2\langle B_2 C_1 C_3 \rangle + 2\langle B_1 C_1 C_4 \rangle + 2\langle B_1 C_2 C_3 \rangle
\end{aligned}$$

这里的  $B_p, C_q$  分别表示  $\zeta_x, \zeta_y$  的第  $p$  阶和  $q$  阶组成波, 系数  $\langle \dots \rangle$  为反映不可约波-波耦合相互作用的量值。

### 2 波面斜率联合概率统计分布

设  $p(\zeta_x, \zeta_y)$  为波面斜率  $\zeta_x, \zeta_y$  的联合概率密度函数, 则  $p(\zeta_x, \zeta_y)$  的特征函数  $\phi(ip, iq)$  可表示为:

$$\phi(ip, iq) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} p(\zeta_x, \zeta_y) \exp(i p \zeta_x + q \zeta_y) d\zeta_x d\zeta_y$$

$$= \sum_{nm} [\langle \zeta_x^n \zeta_y^m \rangle / n! m!] (ip)^n (iq)^m \quad (10)$$

将(5), (6), (7), (8)式代入(10)式, 则有:

$$\phi(ip, iq) = \exp - [(\sigma_p)^2 + (\sigma_q)^2] / 2 \cdot \left[ 1 + \sum_{(n,m) \neq (0,0)}^4 a_{nm} (ip)^n (iq)^m \right] \quad (11)$$

其中相关系数  $a_{nm}$  按(9)式确定。

分布函数  $p(\zeta_x, \zeta_y)$  可由其特征函数  $\phi(ip, iq)$  经 Fourier 逆变换求得:

$$\begin{aligned} p(\zeta_x, \zeta_y) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi(ip, iq) \exp - i(p\zeta_x + q\zeta_y) dp dq \\ &= (1 / 2\pi\sigma_x\sigma_y) \exp - [(\zeta_x / \sigma_x)^2 + (\zeta_y / \sigma_y)^2] / 2 \times \\ &\quad \left[ 1 + \sum_{(n,m) \neq (0,0)}^4 (A_{nm} / n! m!) H_n(\zeta_x / \sigma_x) H_m(\zeta_y / \sigma_y) \right] \end{aligned} \quad (12)$$

其中,  $A_{nm} = (n! m! / \sigma_x^n \sigma_y^m) a_{nm}$ , (12)式即为四阶近似下, 仅当略去四波以上耦合相互作用项时, 给出的非线性随机海浪波面斜率的联合概率统计分布。

### 3 结论

**3.1** 依据 Longuet-Higgins 非线性随机海浪模型, 在高阶近似下, 应用直接求耦合矩的方法, 导出的各向异性波面斜率统计分布为截断的 Gram-Charlier 级数。非线性近似的阶数是与不同量阶的波-波耦合相互作用的类型相联系的, 并最终确定了 Gram-Charlier 级数应截断的项数。

**3.2** 由于考虑到海面风作用方向的影响, 波面斜率统计分布具有非对称性, 仅当忽略波-波耦合相互作用的非线性影响时, 则分布蜕化为正态分布。

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# JOINT STATISTICAL DISTRIBUTION OF SURFACE SLOPES FOR THE HIGHER ORDER NONLINEAR RANDOM SEA WAVES

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**Abstract** The non-Gaussian theory of Sun Fu(1994), a statistical distribution of surface elevation, is extended here to obtain the joint distribution of anisotropy surface slopes for a weakly nonlinear model of random sea waves by applying the method of calculating directly each coupling moment. It is shown that in higher approximation of sea surface, the distribution can be given by the truncated Gram-Charlier series. The types of wave-wave coupling interactions are connected with the order of approximation to nonlinearity of sea surface, which eventually defines the truncated term of the Gram-Charlier series. For each order of approximation, the coefficients in the series are modified comparatively to the corresponding ones for the previous order approximation. The distribution reduces to Gaussian if the effect of nonlinear wave-wave coupling interaction is excluded.

**Key words** Nonlinear statistical distribution    Wave surface slopes    Coupling moment

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